

## Photon Shielding Features of Quarry Tuff

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**Abstract.** Cantera is a quarry tuff widely used in the building industry; in this work the shielding features of cantera were determined. The shielding characteristics were calculated using XCOM and MCNP5 codes for 0.03, 0.07, 0.1, 0.3, 0.662, 1, 2, and 3 MeV photons. With XCOM the mass interaction coefficients, and the total mass attenuation coefficients, were calculated. With the MCNP5 code a transmission experiment was modelled using a point-like source located 42 cm apart from a point-like detector. Between the source and the detector, cantera pieces with different thickness, ranging from 0 to 40 cm were included. The collided and uncollided photon fluence, the Kerma in air and the Ambient dose equivalent were estimated. With the uncollided fluence the linear attenuation coefficients were determined and compared with those calculated with XCOM. The linear attenuation coefficient for 0.662 MeV photons was compared with the coefficient measured with a NaI(Tl)-based  $\gamma$ -ray spectrometer and a <sup>137</sup>Cs source.

### 1 Introduction

Due to the importance for radiation protection the shielding characteristics of man-made or natural occurring materials have been investigated. Thus, different shielding properties of building materials [1-3], ores [4], glasses [5, 6], plastics and polymers [7], Lunar soil samples [8], gemstones, like the Amethyst [9], Boron containing materials [10], alloys [11], nanoparticles of W and polyethylene dispersed in polymer [12], and concrete with different aggregates [13-16] have been reported.

Some of the shielding features include the effective atomic number, the effective electron density, the half value layer, the energy absorption and exposure buildup factors, the linear attenuation coefficient ( $\mu$ ), and the mass attenuation coefficient ( $\mu/\rho$ ) for photons.

In the cases of lunar soils the dose reduction and the protection against Galactic Cosmic Rays have been studied [8].

Another important shielding feature includes the materials behaviour against neutrons, thus Singh, Badiger and El-Khayatt [1] reported the shielding features against neutrons of building materials.

Also designed materials, like the Hormirad<sup>TM</sup>, which is a high-density concrete has been evaluated as neutron shielding using experimental procedures with isotopic neutron sources and through Monte Carlo calculations

[17]. Another designed material, known as NGS-concrete, polymers and standard cements mortars have been studied in their shielding features [18, 19]. In these studies, the half value layer and the ten value layer were evaluated.

The shielding characterization has been carried out through measurements, calculations or combining both procedures.

Tuffs are volcanic rocks made of an ash matrix with grain sizes ranging from fine clay minerals up to silt-sized material, containing lithic, vitric and crystal clasts. Tuffs are formed from pyroclastic flow, surge and fall deposits that combined with the primitive magma features, the deposition and compaction of the emitted material and the post depositional processes concur to find tuffs with different textures, and petrographic and chemical compositions [20].

Volcanic tuffs minerals have been proposed to improve the indoor environments reducing the microbial contaminants and the airborne fungal carcinogens [21].

Quarry tuffs are mostly soft and porous rocks used as building stones and for artwork because can be easily cut and reworked. Volcanic tuff stones are in different colours, they are used as covering materials for insulating and ornamental purposes on the exterior and interior of buildings [22-25].

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In the construction industry the quarry tuff (cantera) is used to coat the interior and exterior wall surfaces of buildings and houses avoiding the use of paint. Its use is part of the Mexico's stone heritage being present in pre-Hispanic, Colonial and modern architecture [26]. A radiometric analysis of this material from Turkey has been reported [24, 25], being in mostly of cases safe.

In Mexico, there are several cities where the cantera is widely used, it is found in prehispanic constructions, cathedrals, and modern buildings and houses. Cantera is used and the main construction material or as decorative item attached to the walls.

In several facilities having x-ray machines or sealed radioactive sources cantera is attached to the walls however, its features to shield x-or- $\gamma$  photons are unknown; therefore in facilities with cantera hosting x-ray units or  $\gamma$ -ray sources the cantera is not accounted for shielding design or to evaluate the shielding effectiveness.

The aim of this work was to estimate the shielding characteristics of cantera using Monte Carlo methods and the XCOM code [27] for 0.03, 0.07, 0.1, 0.3, 0.662, 1, 2, and 3 MeV mono-energetic photons. The relative transmission of Kerma in air,  $K_a$ , and the Ambient dose equivalent,  $H^*(10)$ , was also determined.

For 0.662 MeV photons the calculated  $\mu$  was compared with the  $\mu$  measured using a NaI(Tl)-based  $\gamma$ -ray spectrometer and a  $^{137}\text{Cs}$  source.

## 2 Materials and methods

From the local market a large piece of cantera was purchased and it was cut in 10x10 cm pieces with different thickness, ranging from 1 up to 40 cm. Each piece was weighted to determine the density, being  $1.8 \pm 2\%$   $\text{g cm}^{-3}$ .

From the large piece of cantera several samples were taken in order to measure the chemical composition using Energy Dispersive X-ray Fluorescence. Cantera is mainly composed by  $\text{SiO}_2$ ,  $\text{Al}_2\text{O}_3$ ,  $\text{K}_2\text{O}$ ,  $\text{Na}_2\text{O}$ ,  $\text{Fe}_2\text{O}_3$ , and  $\text{CaO}$ , being in agreement with the chemical composition reported by Celik and Ergul [28]. Due to the total alkali ( $\text{Na}_2\text{O}+\text{K}_2\text{O}$ ) versus silica ( $\text{SiO}_2$ ) composition, this cantera tuff is Rhyolite [28, 29].

### 2.1 Calculations

The chemical composition of cantera was used to calculate the elemental concentration in weight fraction, being: Oxygen (0.4604 w/o), Silicon (0.2476 w/o), Aluminium (0.1022 w/o), Carbon (0.0651 w/o), Iron (0.0509 w/o), Potassium (0.0459 w/o), Sodium (0.0156 w/o), Magnesium (0.0074 w/o), and Calcium (0.0049 w/o). The elemental concentration was used to calculate the shielding features against photons using the XCOM and the Monte Carlo code MCNP5 [30].

#### 2.1.1 XCOM

The XCOM code uses a large data base with Coherent scattering, Compton scattering (Incoherent scattering), Photoelectric absorption, and Pair production cross sections for several elements.

The chemical composition of cantera was input in the XCOM code and the mass interaction coefficients and the total mass attenuation coefficients, for 1 keV to 100 GeV photons, were calculated. The mass attenuation coefficient is calculated using equation 1.

$$\left(\frac{\mu}{\rho}\right) = \sum_i w_i \left(\frac{\mu}{\rho}\right)_i \quad (1)$$

In equation 1,  $\rho$  is the cantera density, and  $w_i$  and  $(\mu/\rho)_i$  are the weight fractions and mass attenuation coefficient of the constituent element  $i$  respectively.

The density of cantera and the total mass attenuation coefficients were used to calculate the linear attenuation coefficient,  $\mu$ , for photons of 0.03, 0.07, 0.10, 0.3, 0.662, 1, 2 and 3 MeV, using equation 2.

$$\mu(E) = \rho \left(\frac{\mu}{\rho}\right)(E) \quad (2)$$

Here,  $\mu(E)$  is the linear attenuation coefficient of cantera for photons of  $E$  MeV,  $\rho$  is the density of cantera, and  $(\mu/\rho)(E)$  is the mass attenuation coefficient of cantera for photons of  $E$  MeV.

#### 2.1.2 MCNP5

In the MCNP5 calculations a model of a transmission experiment was built. Here, a point-like and isotropic photon source was used where the source term was mono-energetic, with photons of 30, 70, 100, and 300 keV in order to cover the x-rays for diagnosis, and 0.662, 1, 2 and 3 MeV to include gamma-ray sources.

A point-like detector was placed 42 cm from the source, between the source and the detector models of the cantera pieces where included. Calculations were carried out for 0, 0.5, 1, 2, ..., 39, 40 cm-thick cantera pieces. The f5 tally was used to estimate the collided (total) and the uncollided photon fluence.

The uncollided photon fluence and the cantera thickness were fitted to the exponential function in order to calculate the linear attenuation coefficients that were compared with the  $\mu$  calculated with the XCOM code.

The MCNP5 calculations also included the Kerma in air,  $K_a$ , and the Ambient dose equivalent,  $H^*(10)$  that were calculated using the ICRP 74 [31] fluence-to-dose conversion coefficients. For each case the amount of histories used in the Monte Carlo calculations was large enough to have an uncertainty less than 5%.

## 2.2 Measurements

Using a  $555 \pm 5\%$  MBq  $^{137}\text{Cs}$  source in a lead shield with a 1 cm-diameter collimator, a transmission experiment was carried out using a narrow geometry array.

The transmitted photon spectrum was measured with gamma-ray spectrometer with a  $7.62 \text{ } \varnothing \times 7.62 \text{ cm}$  NaI(Tl) detector and a multichannel analyser. The detector was sited in a lead shield with 1 cm-diameter collimator. The centres of both collimators were aligned.

The area under the 0.662 MeV photopeak was measured as the cantera pieces with different thickness were placed between the source and the detector. The live-time of measurements was set to get 1% uncertainty in the photopeak net area.

For each cantera thickness the net count rates under the 0.662 photopeak were corrected due to background. Corrected count rates were normalized to the count rate measured without cantera and adjusted to an exponential function using the weighted regression [32]. The weight factors,  $\omega$ , were calculated using equation 3.

$$\omega_i = \frac{I}{sc_i + sx_i} \quad (3)$$

Here,  $sc_i$  is the variance of count rates measured with the  $i$ -th cantera piece,  $sx_i$  is the variance of the  $i$ -th cantera thickness.

## 3 Results

### 3.1 Calculations

#### 3.1.1 XCOM results

The cantera mass interaction coefficients, for 1 keV-100 GeV photons, calculated with the XCOM code, are shown in figure 1.

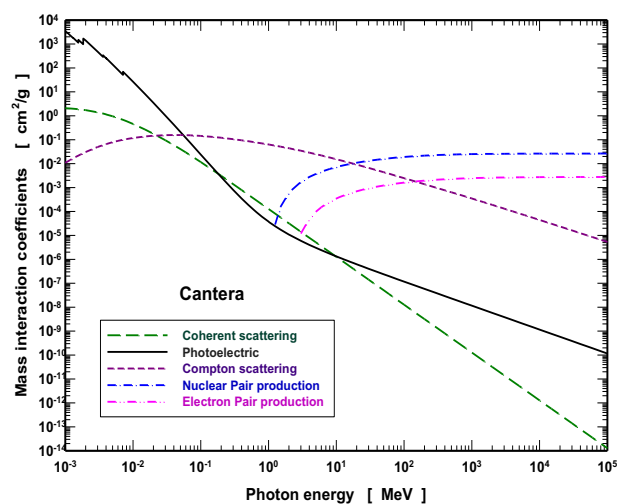


Figure 1. Photon mass interaction coefficients in cantera.

The interaction coefficients include coherent and incoherent (Compton) scattering, photoelectric effect, as well as, nuclear and electron pair production.

In cantera the photoelectric absorption is the main interaction mechanism for photons less than 60 keV, the best shielding performance is for low energy photons. For photons from 60 keV to approximately 20 MeV Compton scattering is the main interaction process with cantera where scattered photon, with lower energy, are produced.

The total mass attenuation coefficients of cantera for  $10^{-3}$  to  $10^5$  MeV photons, are shown in figure 2.

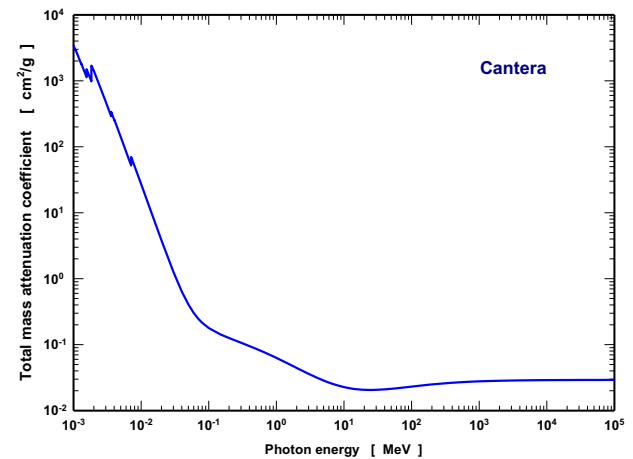


Figure 2. Total mass attenuation coefficients of cantera.

The total mass attenuation coefficients decreases sharply in the low energy region were discrete, small size, peaks appears. This behaviour is similar to the total mass attenuation coefficients reported for building materials [3]. In the low energy region there are seven resonances due to K absorption edge in Na, Mg, Al, Si, K, Ca, and Fe.

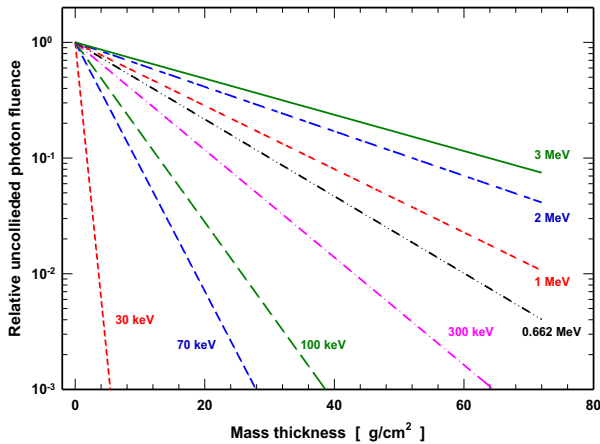
Using the mass attenuation coefficients and the density of cantera the linear attenuation coefficients for 0.03, 0.07, 0.1, 0.3, 0.662, 1, 2 and 3 MeV photons were calculated being included in Table 1.

Table 1. Linear attenuation coefficients of cantera.

Photon energy [MeV]	$\mu_{\text{XCOM}}$ [ $\text{cm}^{-1}$ ]	$\mu_{\text{MCNP5}}$ [ $\text{cm}^{-1}$ ]
0.030	2.2895	2.2920
0.070	0.4457	0.4455
0.100	0.3218	0.3219
0.300	0.1928	0.1926
0.662	0.1380	0.1379
1	0.1137	0.1135
2	0.0798	0.0797
3	0.0650	0.0649

### 3.1.2 Monte Carlo results

In Figure 3 are shown the relative uncollided photon transmission, estimated with the MCNP5 code for monoenergetic photon sources in function of the mass thickness of cantera.



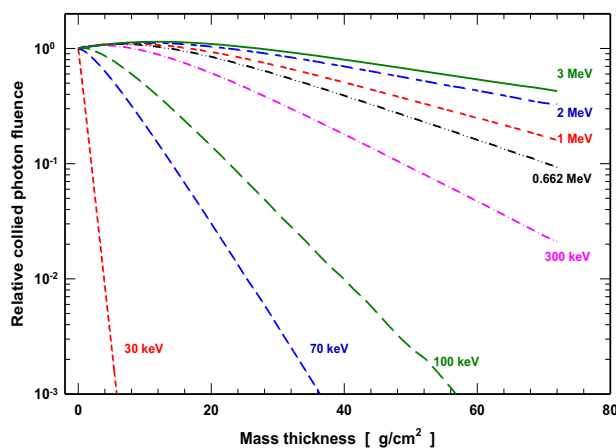
**Figure 3.** Uncollided photon transmission in cantera.

Data from Figure 3 were fitted into exponential functions and the linear attenuation coefficients were calculated.

In Table 1 are shown the  $\mu$  values obtained with the XCOM and MCNP5 codes. There is a good agreement between the  $\mu$  calculated using both codes.

The largest difference, noticed for 3 MeV photons, was 0.27 % being statistically insignificant. The Half value layer of cantera is 0.3, 1.6, 2.1, 3.6, 5.0, 6.1, 8.7, and 10.7 cm for 0.03, 0.07, 0.1, 0.3, 0.662, 1, 2, and 3 MeV photons respectively.

In Figure 4 the relative total photon (collided) transmission in cantera is shown.

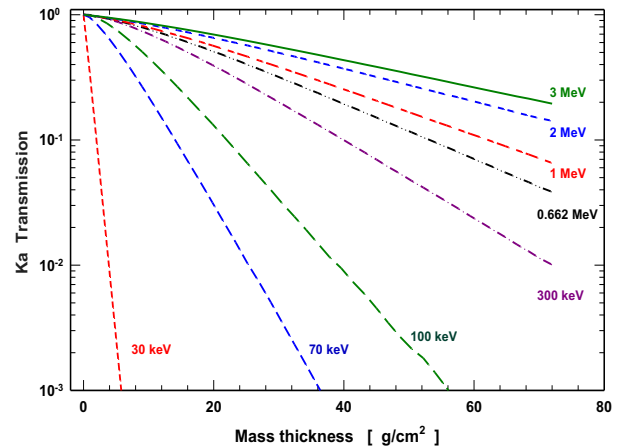


**Figure 4.** Collided photon transmission in cantera.

Total photon fluence includes the uncollided and the scattered photons. The differences between the photon transmission in figures 3 and 4 are due to photon buildup in the detector. The photon buildup is the increase of photon fluence in the detector due to scattered photons by the

cantera that reach the detector. Scattered photons have lower energies.

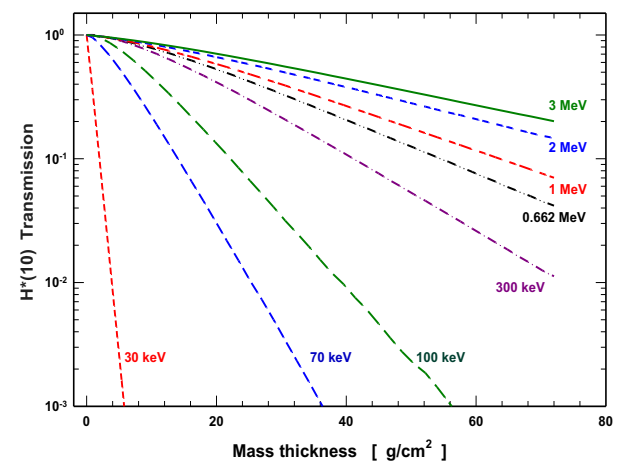
In Figure 5 the Ka transmission due to total photon fluence, in function of cantera mass thickness ( $x/\rho$ ) for mono-energetic photons are presented. This feature is shown in terms of ( $x/\rho$ ) because the density of cantera tuffs varies from 1.80 to 2.37  $\text{g cm}^{-3}$  [22].



**Figure 5.** Kerma in air transmission in cantera.

The Ambient dose equivalent transmission of mono-energetic photons in cantera is shown in Figure 6.

Despite the fluence-to- $H^*(10)$  conversion coefficients are larger than the fluence-to-Ka conversion coefficients for photons from 0.03 to 3 MeV [31], for the same thickness of cantera the  $H^*(10)$  transmission is slightly larger than Ka transmission. The probable explanation is because  $H^*(10)$  and the Ka values were calculated using collided photons (uncollided and scattered) that reach the detector, therefore some scattered photons has energies where fluence-to-Ka conversion coefficients are larger than fluence-to- $H^*(10)$  conversion coefficients.



**Figure 6.** Ambient dose equivalent transmission in cantera.

## 3.2 Measurements

The linear attenuation coefficient of cantera for 0.662 MeV photons from  $^{137}\text{Cs}$  source is  $0.14 \pm 0.01 \text{ cm}^{-1}$ . This

value was obtained through measurements carried out with the NaI(Tl)  $\gamma$ -ray spectrometer. The measured  $\mu$  is in agreement with the  $\mu$ -values for 0.662 MeV photons calculated with XCOM and MCNP5 codes, shown in Table 1.

The  $\mu/\rho$  for 0.662 MeV is  $0.076 \text{ cm}^2 \text{ g}^{-1}$  being similar to the value shown in building materials [2] and is larger than  $\mu/\rho$  of aluminium ( $7.466 \times 10^{-2} \text{ cm}^2 \text{ g}^{-1}$ ) and smaller than  $\mu/\rho$  of lead ( $0.1101 \text{ cm}^2 \text{ g}^{-1}$ ) [33].

## 4 Conclusions

The cantera is widely used in the construction industry however, it's shielding features against x and  $\gamma$ -rays were unknown. Through calculations with XCOM and MCNP5 codes the shielding features of cantera have been estimated for 0.03, 0.07, 0.100, 0.300, 0.662, 1, 2, and 3 MeV monoenergetic photons. Using a  $^{137}\text{Cs}$  and a NaI(Tl)  $\gamma$ -ray spectrometer the linear attenuation coefficient for 0.662 MeV photons was also measured.

Due to total alkali and silica content in this quarry tuff is defined as rhyolite having a density of  $1.8 \text{ g cm}^{-3}$ .

Linear attenuation coefficient of cantera varies from approximately  $2.290$  to  $0.065 \text{ cm}^{-1}$  for 30 keV to 3 MeV photons respectively.

For 30 keV photons the half value layer of cantera is 0.3 cm and 10.7 cm for 3 MeV  $\gamma$ -rays.

Photons with energy less than 60 keV are effectively shielded because the main interaction with cantera is through photoelectric interaction. Thus, the cantera is a good shielding material for rooms with x-ray units working below 60 kV. This x-ray equipment includes radio diagnostic x-rays for mammography or some few used for dental radiographies.

For photons with energy between 60 keV and 10 MeV the main interaction in cantera is Compton scattering. Here, the energy of the incoming photon is reduced, but scattered photons contribute with the photon and the dose build up that should be evaluated. Nevertheless in rooms with x-ray units working above 60 kV, like dental, diagnosis, interventional x-rays, or those with  $\gamma$ -ray sources, the cantera in the walls will help to shield photons and to reduce the dose.

For the same thickness of cantera the Ka and H\*(10) transmission is almost the same.

A limitation of this work was the use of monoenergetic photons to represent x-rays, because x-ray tubes produce photons with a wide and continuous energy distribution together with few discrete peaks where just quite few photons have the largest energy. However, if the information here presented is used to calculate a shielding made with cantera will be conservative because

the photon mean energy of actual x-rays is smaller than the energy of photons here used.

Another limitation of this work was that photons and dose buildup factors for cantera were not included.

## References

1. V. P. Singh, N. M. Badiger and A. M. El-Khayatt, *Radiat. Eff. Def. Solids* **169**, 547 (2014)
2. K. S. Mann, B. Kaur, G. S. Sidhu, A. Kumar, *Radiat. Phys. Chem.* **87**, 16 (2013)
3. Ch. Singh, T. Singh, A. Kumar, G. S. Mudahar, *Ann. Nucl. Energy* **31**, 1199 (2004)
4. B. Oto, N. Yildiz, F. Akdemir, E. Kavaz, *Prog. Nucl. Energy* **85**, 391 (2015)
5. S. Kaewjang, U. Maghanemi, S. Kothan, H. J. Kim, P. Limkitjaroenporn and J. Kaewkhao, *Nucl. Eng. Des.* **280**, 21 (2014)
6. Y. Gedikoglu, G. Gedikoglu, G. Berkin, T. Ceyhan, M. A. Altinoz, *Toxicol. Ind. Health* **28**, 708 (2012)
7. K. S. Mann, A. Rani and M. S. Heer, *Rad. Phys. Chem.* **106**, 247 (2015)
8. J. Miller, L. Taylor, C. Zeitlin, L. Heilbronn, S. Guetersloh, M. DiGiuseppe, Y. Iwata, T. Murakami, *Radiat. Meas.* **44**, 163 (2009)
9. T. Korkut, H. Korkut, A. Karabulut, G. Budak, *Ann. Nucl. Energy* **38**, 56 (2011)
10. V. P. Singh, N. M. Badiger, *Vacuum*, **113**, 24 (2015)
11. D. N. Hopkins, M. Maqbool, M. S. Islam, *Radiol. Phys. Technol.* **5**, 229 (2012)
12. J. Kim, D. Seo, B. Ch. Lee, Y. S. Seo, W. H. Miller, *Adv. Engr. Mat.* **16**, 1083 (2014)
13. E.-S. A. Waly, M. A. Bourham, *Ann. Nucl. Energy* **85**, 306 (2015)
14. I. I. Bashter, *Ann. Nucl. Energy* **24**, 1389 (1997)
15. I. Akkurt, H. Akyildirim, B. Mavi, S. Kilincarslan, C. Basygit, *Ann. Nucl. Energy* **37**, 910 (2010)
16. A. M. El-Khayatt, *Ann. Nucl. Energy* **37**, 991 (2010)
17. E. Gallego, A. Lorente and H. R. Vega-Carrillo, *Nucl. Tech.* **162**, 399 (2009)
18. T. Piotrowski, D. B. Tefelski, J. J. Sokolowska, B. Jaworska, *Acta Phys. Pol. A* **128**, B-9 (2015)
19. T. Piotrowski, D. B. Tefelski, J. Skubalski, A. Zak, *Acta Phys. Pol. A* **128**, B-14 (2015)
20. P. Ciccioli, P. Plescia, D. Capitani, *J. Phys. Chem. C* **114**, 9328 (2010)
21. P. Yasaka, N. Pattanaboonmee, H. J. Kim, P. Limkitjaroenpron, J. Kaewkaho, *Ann. Nucl. Energy* **68**, 4 (2014)
22. R. López-Doncel, W. Wedekind, T. Leiser, S. Molina-Maldonado, A. Velasco-Sánchez, R. Dohrmann, A. Kral, A. Wittenborn, A. Aguillón-Robles, S. Siegesmund, *Environ. Earth Sci.* **75**, 212 (2016)
23. W. Wedekind, R. López-Doncel, R. Dohrmann, M. Kocher, S. Siegesmund, *Environ. Earth Sci.* **69**, 1203 (2013)
24. M. Degerlier, *Radioprotection* **48**, 215 (2013)
25. S. Turhan, E. Atici, A. Varinlioglu, *Radioprotection* **50**, 273 (2015)

26. N. A. Pérez, E. Lima, P. Bosch, J. Méndez-Vivar, J. Cult. Herit. **15**, 352 (2014)
27. M. J. Berger, J. H. Hubbell, S. M. Seltzer, J. Chang, J. S. Coursey, R. Sukumar, D. S. Zucker, K. Olsen, XCOM: Photon Cross Sections Database, version 1.5, [Online]. <<http://physics.nist.gov/pml/data/xcom/data/xcom/index.cfm>>. Gaithersburg, MD, USA.
28. M. Y. Celik, A. Ergul, Environ. Earth Sci. **74**, 3223 (2015)
29. M. J. Le Bas, R. W. Le Maitre, A. R. Woolley, Miner. Pet. **46**, 1 (1992)
30. X-5 Monte Carlo Team, Report LA-UR-03-1987 Los Alamos National Laboratory (2003)
31. ICRP74, Ann. ICRP **26**, 179 (1996)
32. H. R. Vega-Carrillo, Rev. Mex. Fis. **35**, 597 (1989)
33. T. E. Johnson. B. K. Birky, *Health Physics and Radiological Health* (Lippincott Williams & Wilkins, 2012)